



# Dark matter and entropy dilution

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1205.0844

*ν* Dark workshop  
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# Dark Matter

## stability

# Discrete symmetries

— MSSM & R-parity

— extended Higgs  
Inert Higgs doublet

$Z_2$  symmetries

talk by Steffen

## Discrete symmetries

— MSSM & R-parity

— extended Higgs  
Inert Higgs doublet

$Z_2$  symmetries

talk by Steffen

## Small couplings and or large scales

— small couplings (Dirac Yukawa)

— large scales (extended gauge symmetry)

— i.e. small masses (e.g. axion)

$$\Gamma_\mu \propto m_\mu^5 / M_W^4$$

$$m_{DM} \ll \Lambda$$

talk by Redondo

# Neutrino mass vs. Dark Matter

 new physics needed for neutrino mass

talk by Schwetz

# Neutrino mass vs. Dark Matter

- new physics needed for neutrino mass, same for DM?

# Neutrino mass vs. Dark Matter

— small couplings (Dirac Yukawa)  $\nu_S$

$$\tau_{\nu_S} \simeq \tau_\mu \left( \frac{m_\mu}{m_S} \right)^5 \frac{m_S}{m_\nu}$$

$$m_{\nu_S} \lesssim 20 \text{ keV}$$

$\nu$ MSM

Dodelson, Widrow '93, ...  
can't do justice to the field

talk by Drewes

# Neutrino mass vs. Dark Matter

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Dodelson, Widrow '93, ...  
can't do justice to the field

$$m_{\nu_S} \lesssim 20 \text{ keV}$$

large scales and small masses  $N_i$

Left-Right

no direct channels below  $m_\pi$ , only mixing

Bezrukov,  
Hettmansperger,  
Lindner '09

$$m_{N_1} \simeq \text{keV} \quad M_{W_R} \gtrsim 14 \text{ TeV}$$

a possible window  $M_{W_R} \in (4, 7) \text{ TeV}$

MN, Senjanović, Zhang '12

DM predicts mass and flavor of  $N_i$



# Left-Right the minimal model $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$

→ understanding the SM asymmetry...

Pati, Salam '74  
Mohapatra, Pati '75

...through spontaneous breaking

Senjanović, Mohapatra '75

$$W_L \quad Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$$
$$L_L = \begin{pmatrix} e_L \\ \nu_L \end{pmatrix}$$

$$Q_R = \begin{pmatrix} u_R \\ d_R \end{pmatrix}$$
$$L_R = \begin{pmatrix} e_R \\ \nu_R \end{pmatrix} \quad W_R$$

# Left-Right the minimal model

understanding the SM asymmetry...

Pati, Salam '74  
Mohapatra, Pati '75

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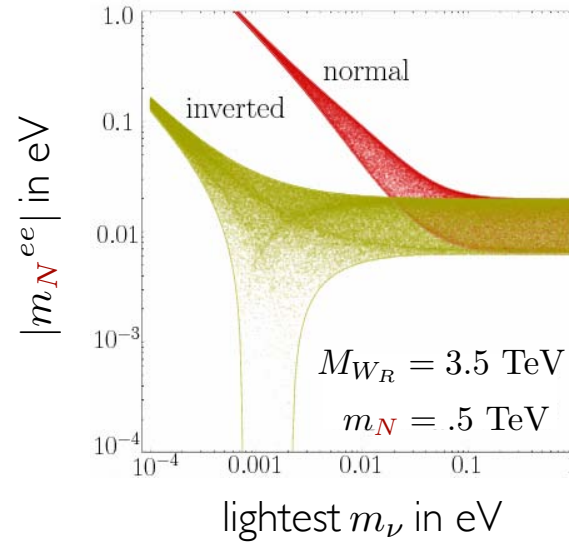
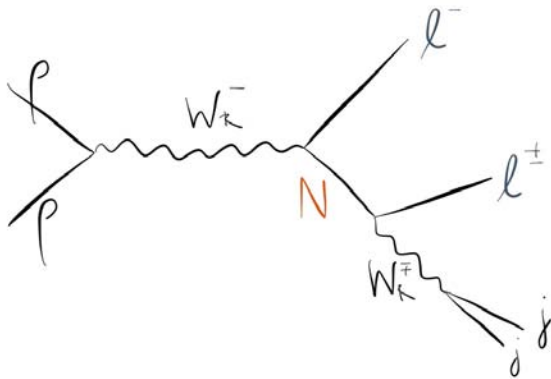
$$\begin{array}{l} W_L \\ Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix} \\ L_L = \begin{pmatrix} e_L \\ \nu_L \end{pmatrix} \end{array} \quad \begin{array}{l} \Phi(2, 2, 0) \\ \Delta_L(3, 1, 2) \quad \Delta_R(1, 3, 2) \end{array} \quad \begin{array}{l} Q_R = \begin{pmatrix} u_R \\ d_R \end{pmatrix} \\ L_R = \begin{pmatrix} e_R \\ \nu_R \end{pmatrix} \\ W_R \end{array}$$

seesaw

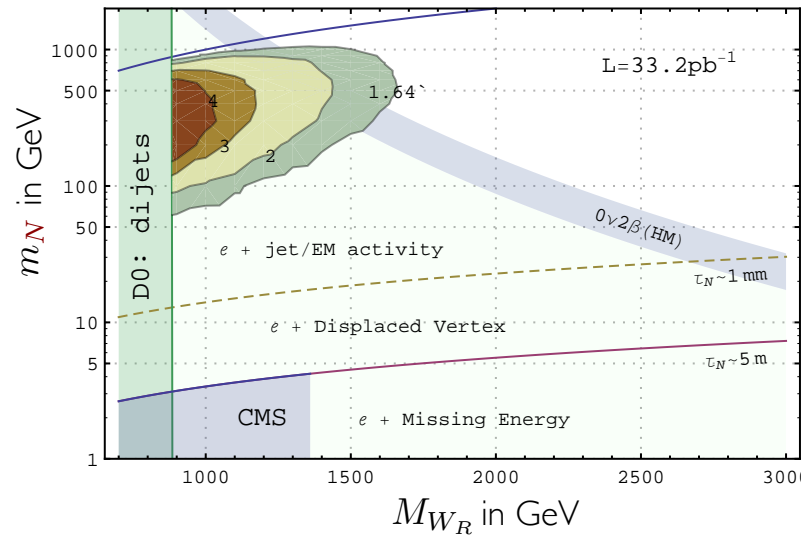
Minkowski '79  
Senjanović, Mohapatra '79

rich phenomenology: LNV, LFV at colliders,  $0\nu 2\beta$

# Left-Right $0\nu 2\beta$ vs. LHC



Senjanović, Mohapatra '82  
Tello, MN, Nesti,  
Senjanović, Vissani '10

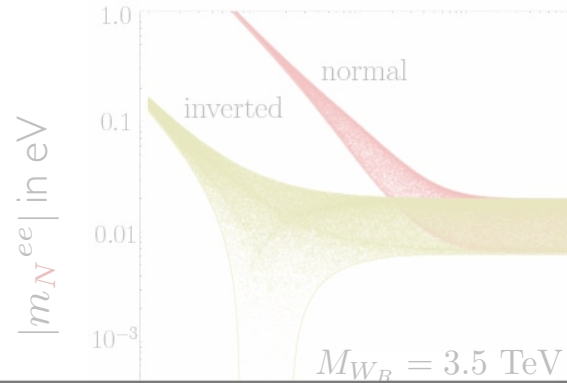


Keung, Senjanović '83  
MN, Nesti, Senjanović,  
Zhang '11

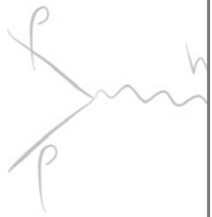
mass origin in Higgs decays  $h \rightarrow N_i N_i$

Maiezza, MN, Nesti, '15

# Left-Right $0\nu 2\beta$ vs. LHC

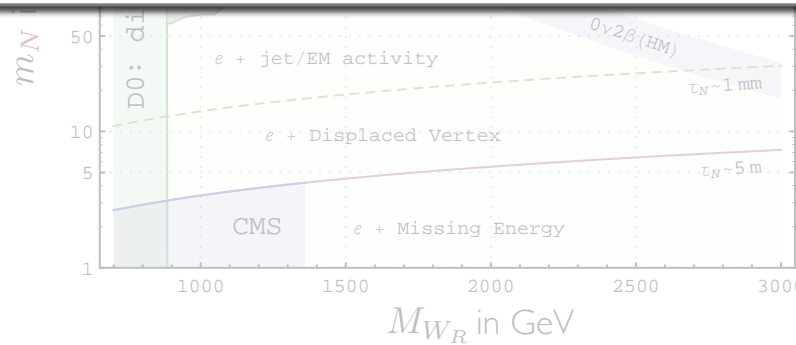


Senjanović, Mohapatra '82  
Tello, MN, Nesti,  
Senjanović, Vissani '10



# Dark Matter?

Senjanović '83  
Senjanović,  
Lang '11



mass origin in Higgs decays  $h \rightarrow N_i N_i$

Maiezza, MN, Nesti, '15

Dark Matter  
(over)production

## Left-Right freeze-out part I

- gauge interactions keep thermal equilibrium
- suppressed rate compared to SM neutrinos

$$\Gamma_N \simeq \Gamma_\nu \left( \frac{M_W}{M_{W_R}} \right)^4 \simeq G_F^2 \left( \frac{M_W}{M_{W_R}} \right)^4 T^5$$

thermal freeze-out

$$\left( \frac{M_W}{M_{W_R}} \right)^4 T^5 \simeq \sqrt{G} T^2$$

around the QCD  
phase transition

$$T_f \simeq 400 \text{ MeV} \left( \frac{g_*(T_f)}{70} \right)^{1/6} \left( \frac{M_{W_R}}{5 \text{ TeV}} \right)^{4/3}$$

## Left-Right over-abundance

- freeze-out relativistic, no Boltzmann suppression

$$T_f > m_N \sim \text{keV}$$

- freeze-out relativistic, no Boltzmann suppression

$$Y_N \equiv \frac{n_N}{s} \simeq \frac{135 \zeta(3)}{4\pi^4 g_*(T_f)}$$

- $Y_N$  conserved, evolve to present times

$$\begin{aligned}\Omega_{N_1} &= Y_{N_1} m_{N_1} \frac{s}{\rho_c} \\ &= 3.3 \left( \frac{m_{N_1}}{\text{keV}} \right) \left( \frac{70}{g_*(T_f)} \right) \\ &\simeq 13 \Omega_{DM}\end{aligned}$$

- many species, not in LRSM

$$g_* < 1000$$

- ~~$m_{N_1} < 0.08 \text{ keV}$~~  astrophysics, ...

# Dark Matter dilution



# Decaying particles in the early universe

Scherrer, Turner '85

- massive late decaying particle
- decays heat up radiation, dilutes the rest  $\equiv$  **diluter**
- applied to sterile neutrinos

Asaka, Shaposhnikov, Kusenko '06

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## Left-Right candidates

- Higgs sector heavy & unstable
- next-to-lightest  $N_{2,3}$  the only remaining candidates

## Diluting with $N_2$ part I

Scherrer, Turner '85

— sudden decay approximation (works for  $\rho_m \gg \rho_\gamma$ )

$$t \leq \tau_2$$

matter dominance

$$H^2 = \Gamma_2^2 = \frac{8\pi}{3M_P^2} (Y_2 m_{N_2}) \frac{2\pi^2 g_*}{45} T_<^3$$

$$T_< = 0.65 \left( \frac{\Gamma_2^2 M_P^2}{Y_2 m_2 g_*} \right)^{1/3}$$

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radiation dominance

$$H = \Gamma_2 = 1.66 \frac{g_*^{1/2}}{M_P} T_{>}^2$$

$$T_{>} \simeq 0.78 g_*^{-1/4} \sqrt{\Gamma_2 M_P}$$
$$\simeq 1.22 \text{ MeV} \sqrt{\frac{1 \text{ sec}}{\tau_2}}$$

needed  
for BBN

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Dilution factor  $\mathcal{S} = \left( \frac{T_{>}}{T_{<}} \right)^3 \simeq 1.7 \frac{T(\tau_2)^{1/4} Y_2 m_2}{\sqrt{\Gamma_2 M_P}}$

$$\mathcal{S} \propto m_2 \sqrt{\tau_2}$$

needed  
for BBN

## Diluting with $N_2$ part I

— how much dilution?

$$\Omega_{N_1} \simeq .228 \left( \frac{m_{N_1}}{\text{keV}} \right) \left( \frac{1.85 \text{ GeV}}{m_{N_2}} \right) \left( \frac{\text{sec}}{\tau_2} \right) \left( \frac{g_*(T_{f2})}{g_*(T_{f1})} \right)$$

Boltzmann  
suppressed

$$T_{f1} \simeq 400 \text{ MeV}$$

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1 sec lifetime near the  $m_\pi + m_\ell$  threshold,  $\tau$  disfavored

$$\tau(N_2 \rightarrow \ell\pi) = \text{sec} \left( \frac{250 \text{ MeV}}{m_{N_2}} \right)^3 \left( \frac{M_{W_R}}{5 \text{ TeV}} \right)^4 \left( \frac{0.002}{f(x_\ell, x_\pi)} \right)^3 \quad \ell = e, \mu$$

# Diluting with $N_2$ part I

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freeze-out flavor dependent

$$m_{N_1} \sim \text{keV}$$

$$V_{Re1} \simeq 0$$

SN cooling

Raffelt, Seckel '88  
Barbieri, Mohapatra '89

'turn off'  $W_R$

$$V_{R\tau 1} \simeq 1$$

$$T_{f1} > T_{f2}$$



# Mass spectrum of DM in Left-Right

— mass and flavor fixed

DM couples to  $\tau$

$$V_R \simeq \begin{pmatrix} 0 & 0 & \textcircled{1} \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

$$m_{N_1} \simeq \text{keV}$$

# Mass spectrum of DM in Left-Right

mass and flavor fixed

$$V_R \simeq \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \begin{aligned} m_{N_2} &\simeq m_\pi + m_\mu \\ m_{N_3} &\simeq m_\pi + m_e \end{aligned}$$

Diluters couple to  $e$  and  $\mu$

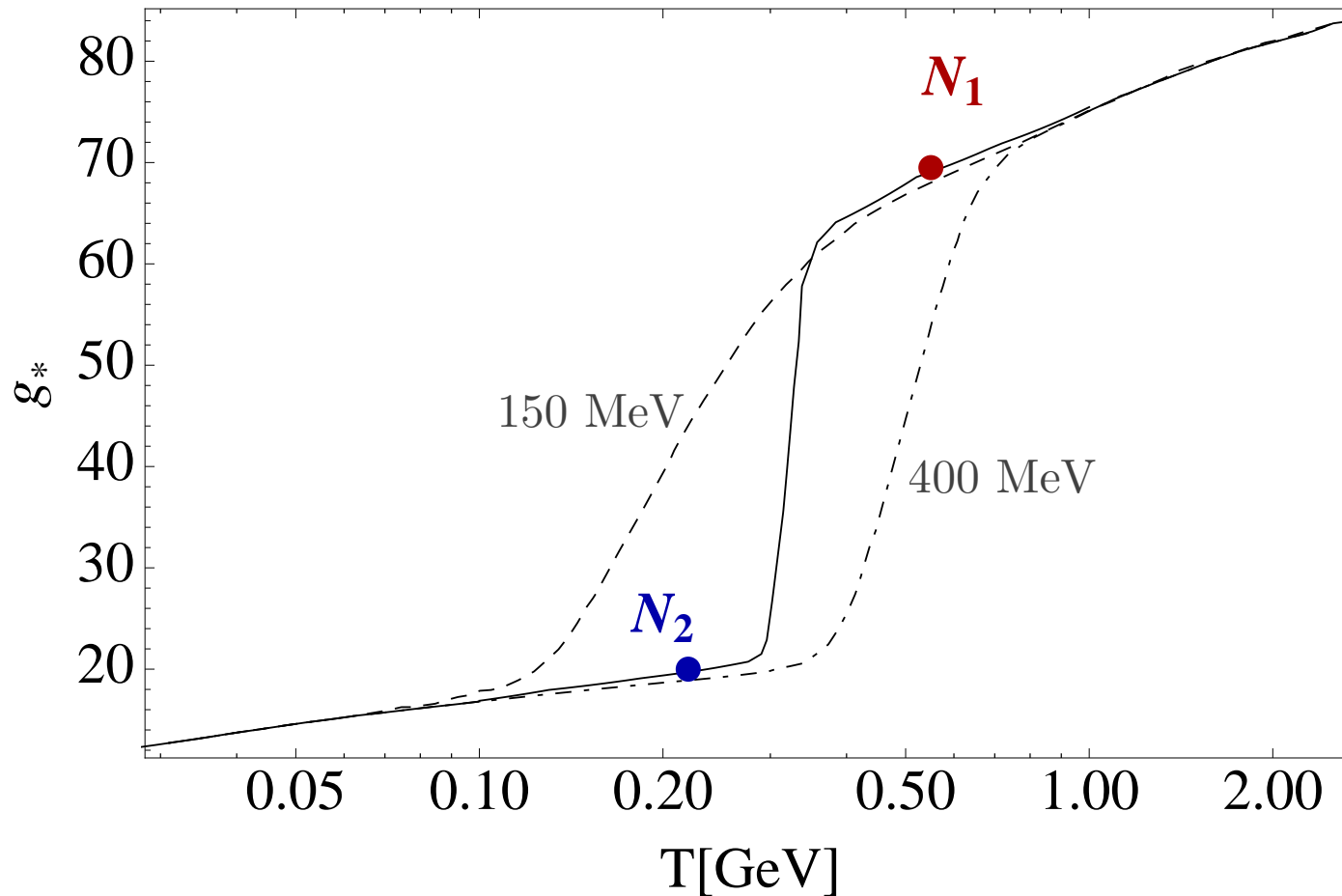
# Dark Matter

## freeze-out pt. 2

# Left-Right freeze-out part II

— drop in  $g_*$

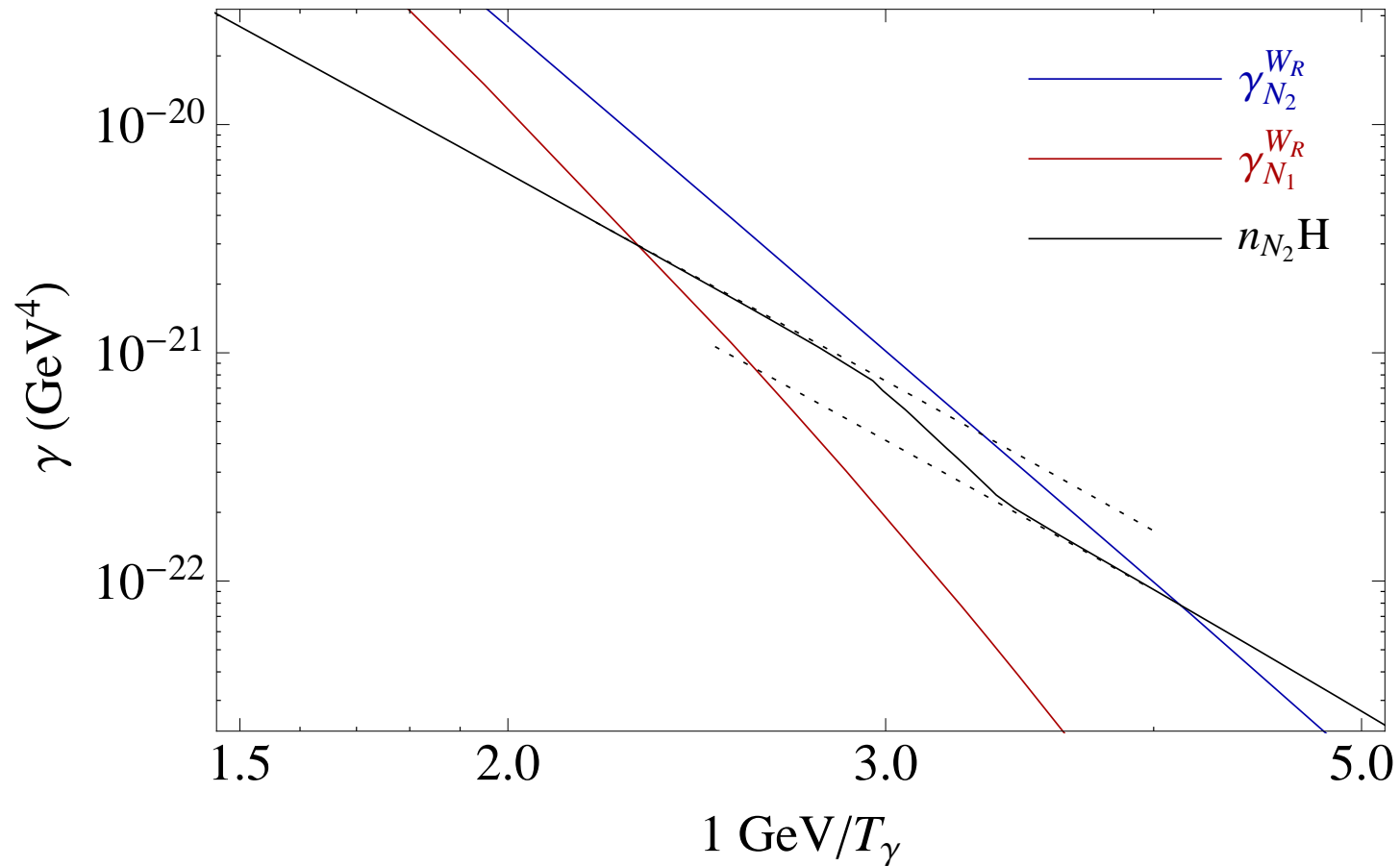
$$\frac{Y_{N_2}}{Y_{N_1}} \simeq \frac{g_*(T_{f1})}{g_*(T_{f2})} \sim 3 - 4$$



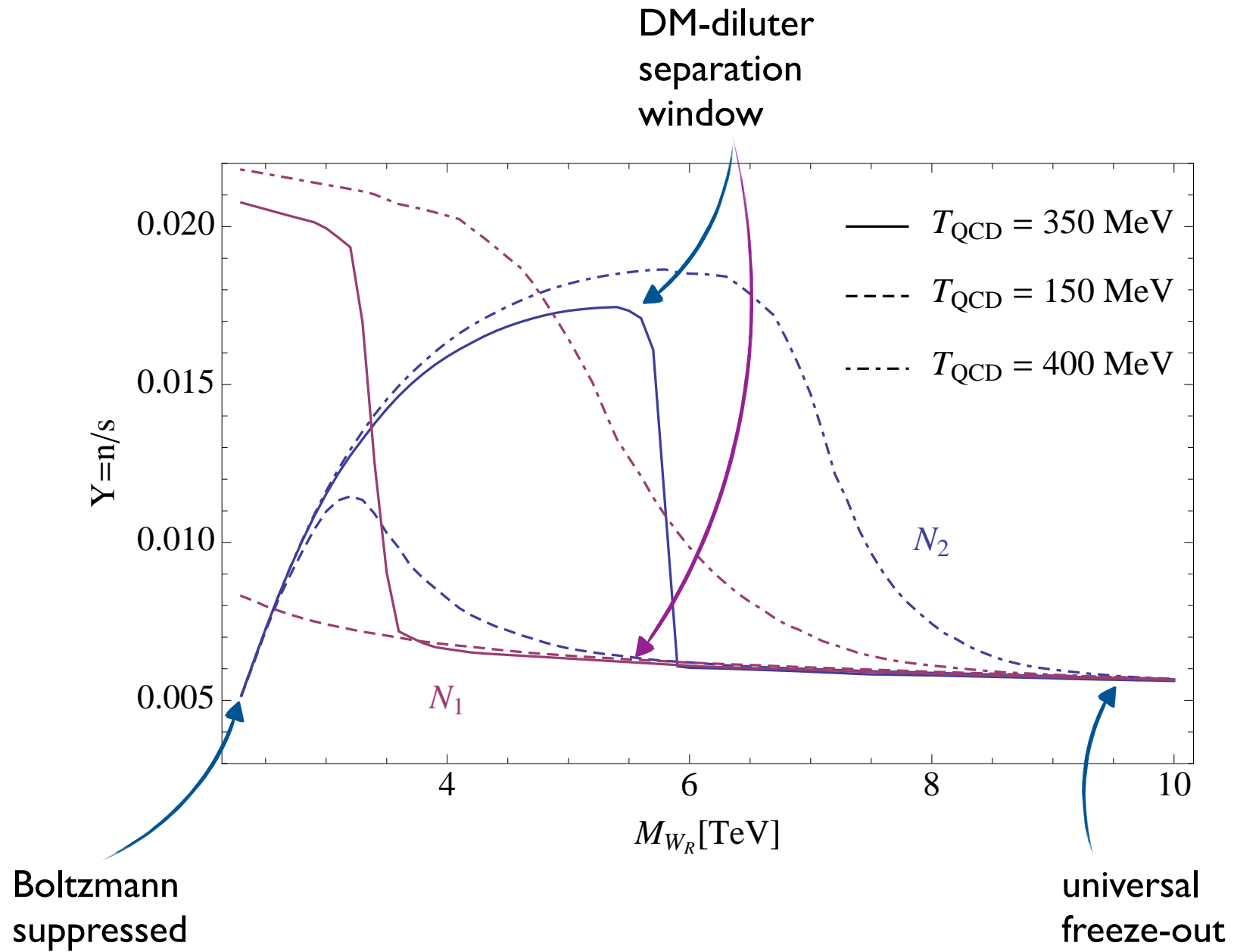
# Left-Right freeze-out part II

pair  $N_i N_i \rightarrow \ell^+ \ell^-, u\bar{u}, d\bar{d}$

single  $N_i \ell^- \rightarrow \bar{u}d, N_i u \rightarrow \ell^+ d, N_i d \rightarrow \bar{\ell}^+ u$



# Left-Right freeze-out part II

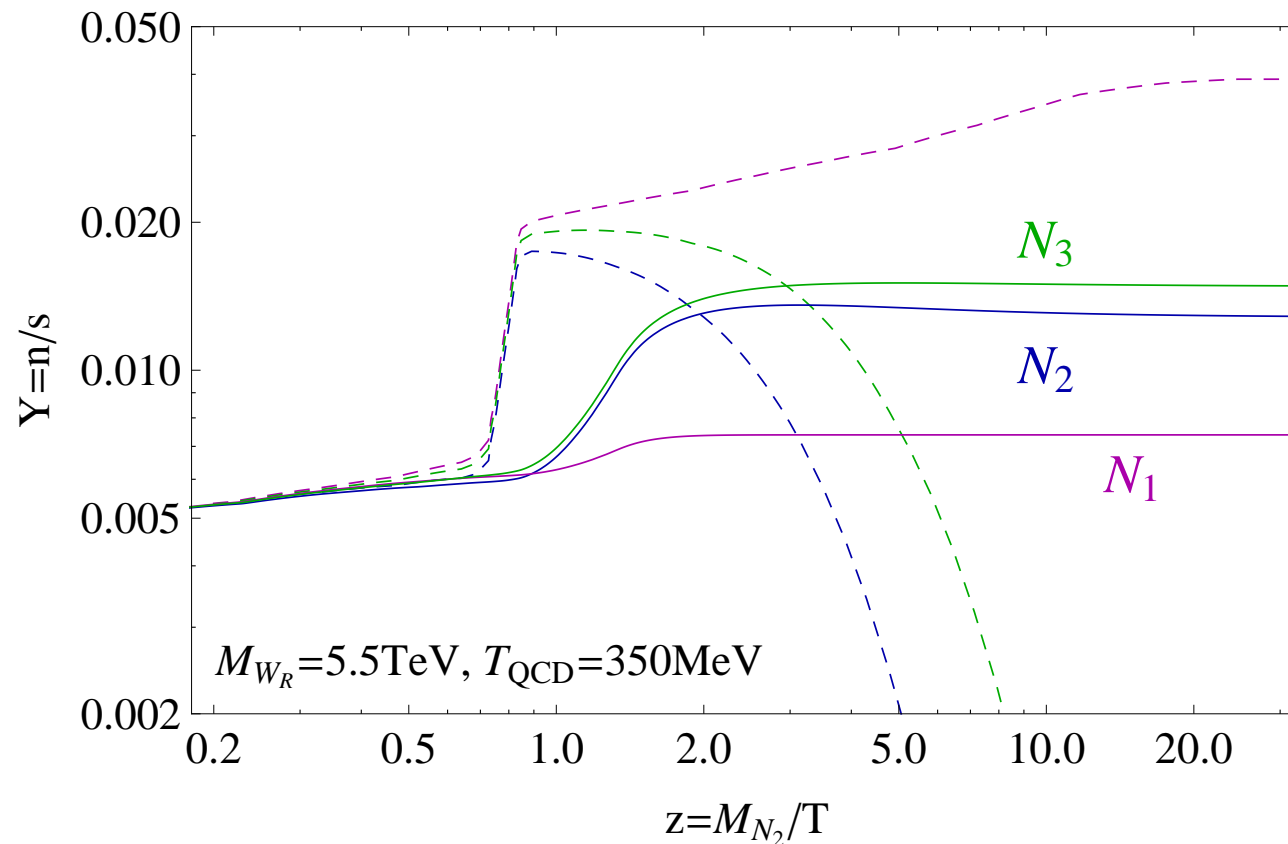


# Left-Right freeze-out part II

coupled Boltzmann equations

$i = 1, 2, 3$

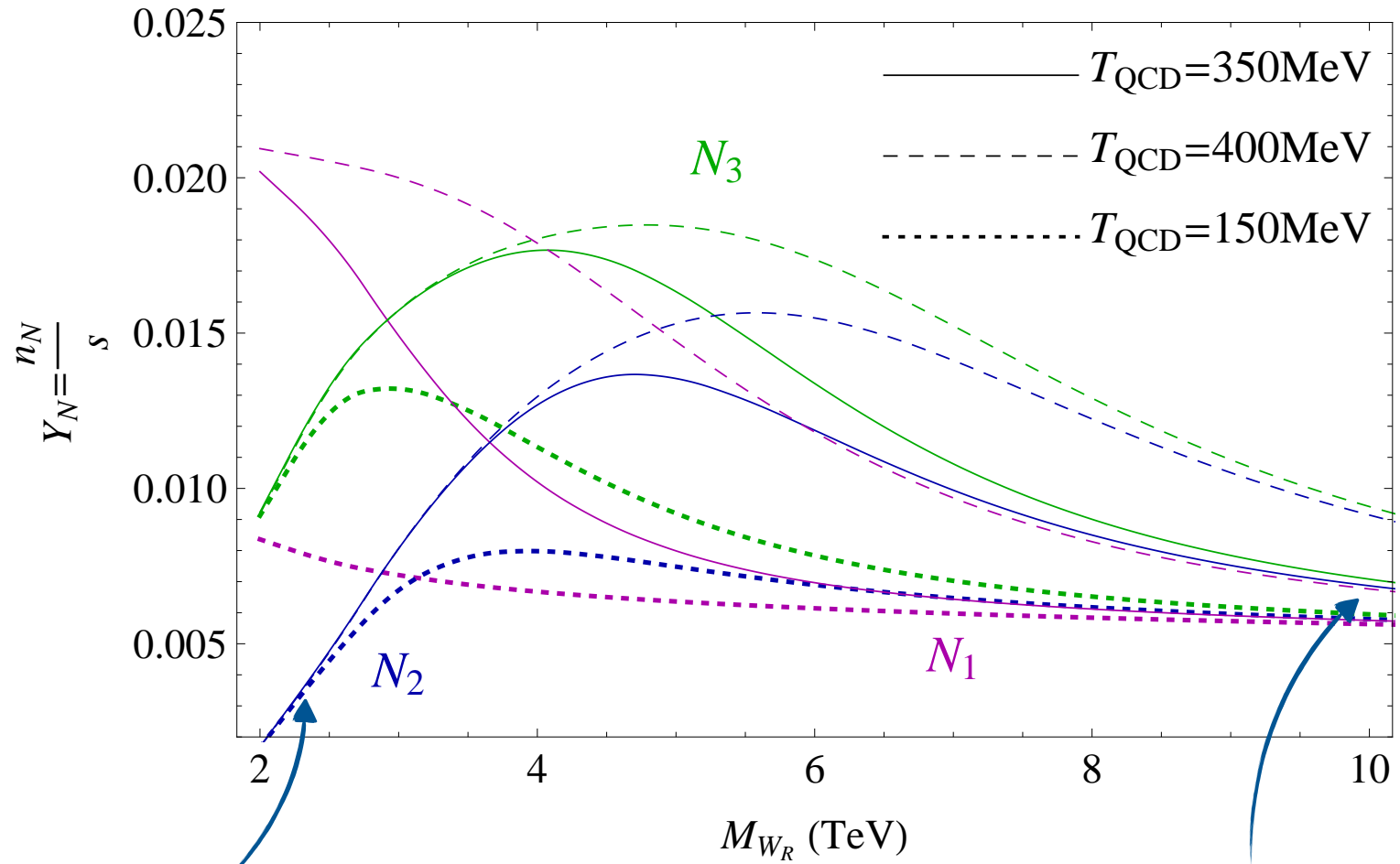
$$sH z \frac{dY_{N_i}}{dz} = - \left( \frac{Y_{N_i}}{Y_{N_i}^{eq}} - 1 \right) \gamma_{N_i}^{W_R} - \left( \left( \frac{Y_{N_i}}{Y_{N_i}^{eq}} \right)^2 - 1 \right) \left( \gamma_{N_i N_i}^{W_R} + \gamma_{N_i N_i}^{Z_{LR}} \right)$$



$$Y_{N_{2,3}} > Y_{N_1}$$

# Left-Right freeze-out part II

— final yields



Boltzmann suppressed

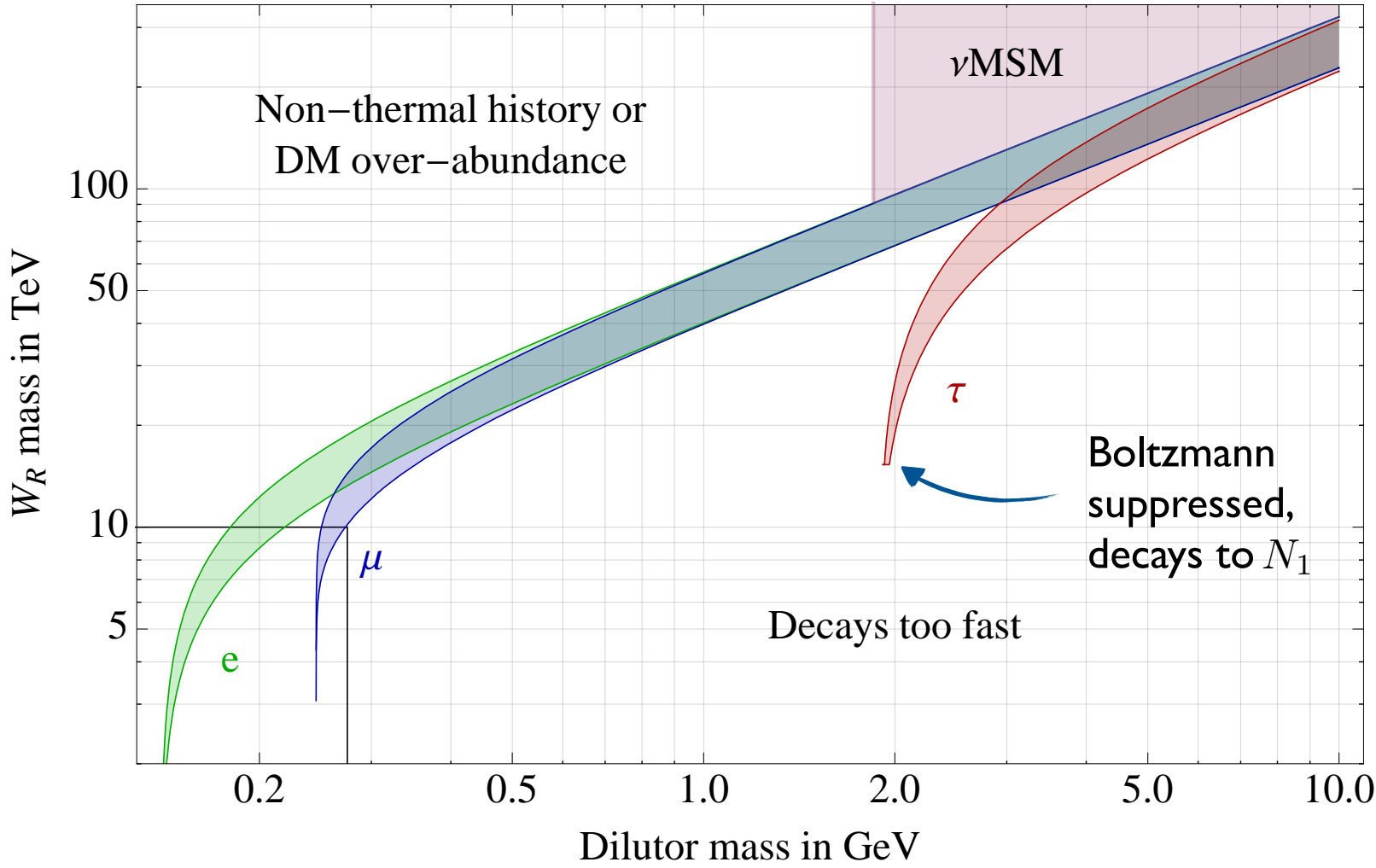
universal freeze-out



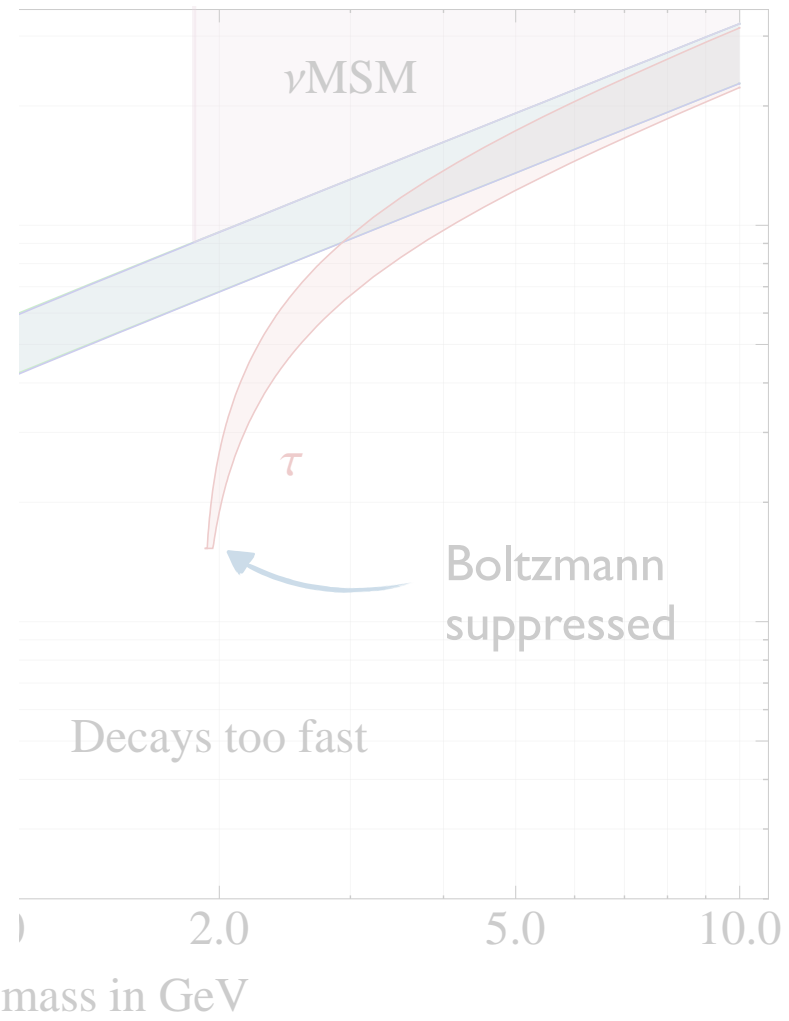
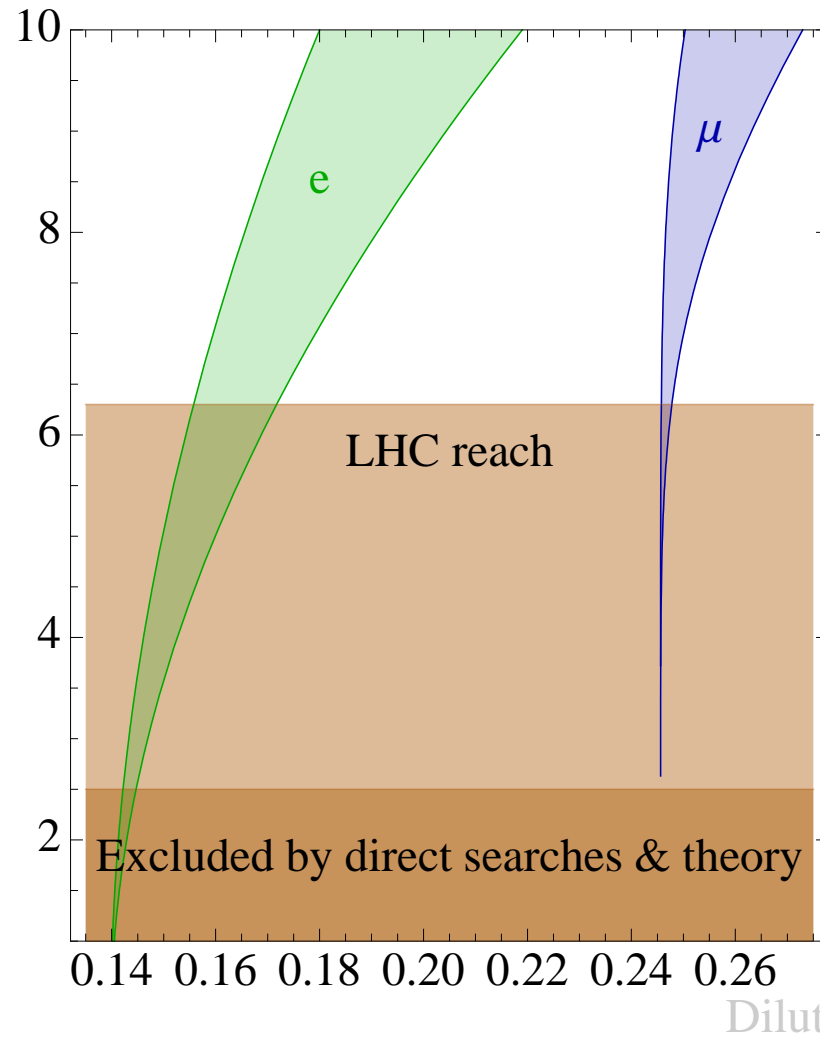
# Dark Matter

## dilution pt. 2

# Dilution phase space



# Dilution phase space



## Dilution part II

← coupled energy density evolution

radiation

$$\frac{d\rho_\gamma}{dt} + 4H\rho_\gamma = \Gamma_2\rho_2 + \Gamma_3\rho_3$$

matter

$$\frac{d\rho_{N_1}}{dt} + 4H\rho_{N_1} = 0$$

$$\frac{d\rho_{N_2}}{dt} + 4H\rho_{N_2} = -\Gamma_2\rho_2$$

$$\frac{d\rho_{N_3}}{dt} + 4H\rho_{N_3} = -\Gamma_3\rho_3$$

# Dilution part II

coupled energy density evolution

radiation

$$\frac{d\rho_\gamma}{dt} + 4H\rho_\gamma = \Gamma_2\rho_2 + \Gamma_3\rho_3$$

matter

$$\frac{d\rho_{N_1}}{dt} + 4H\rho_{N_1} = 0$$

$$\frac{d\rho_{N_2}}{dt} + 4H\rho_{N_2} = -\Gamma_2\rho_2$$

$$\frac{d\rho_{N_3}}{dt} + 4H\rho_{N_3} = -\Gamma_3\rho_3$$

freeze-out

match

$$z_m = 30$$

$$T_{\gamma,m} \simeq 15 \text{ MeV}$$

$$t_m = 1/(2H)$$

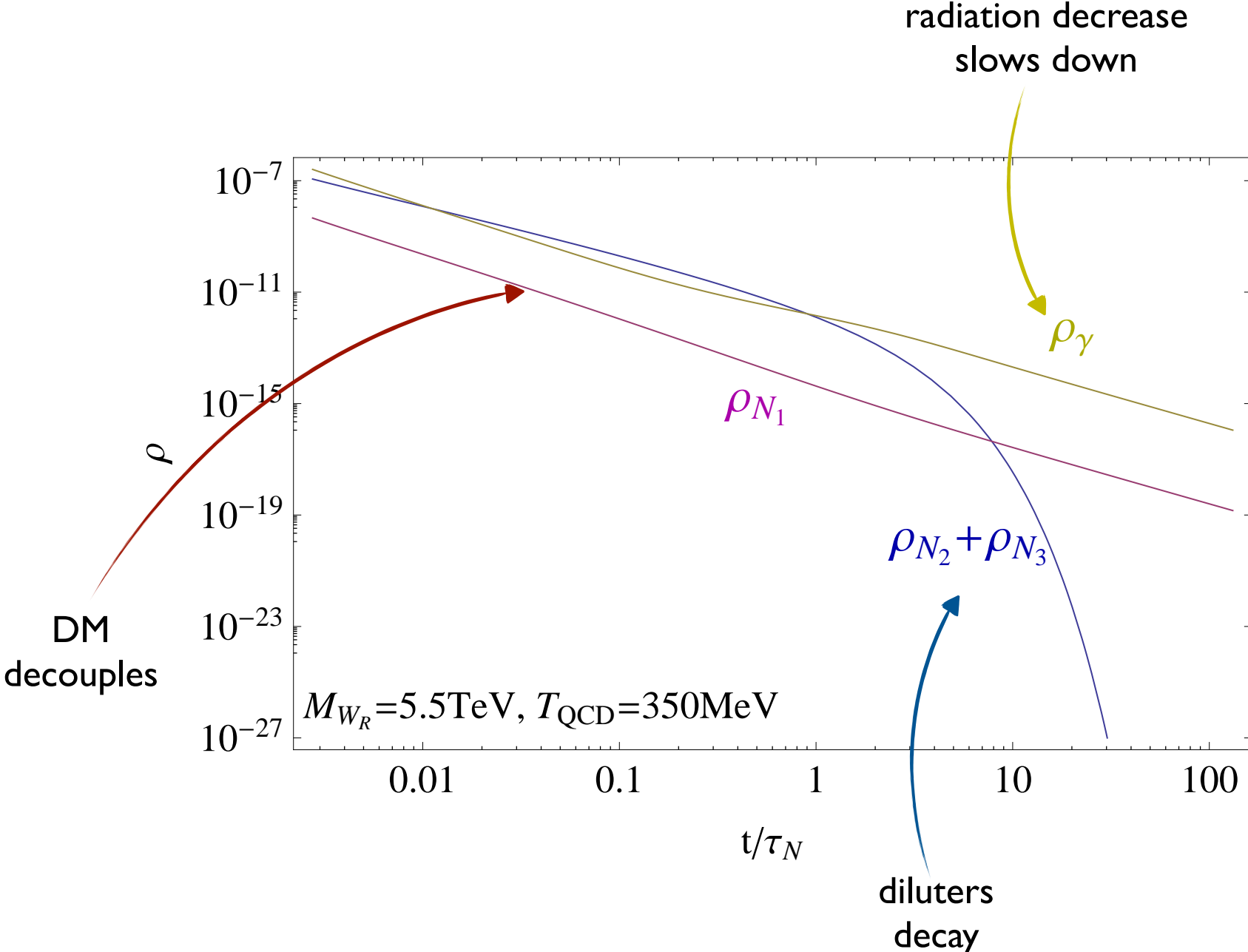
$$\rho_\gamma(t_m) = \frac{\pi^2}{30} g_*(T_{\gamma,m}) T_{\gamma,m}^4$$

$$\rho_{N_1}(t_m) = \frac{7}{4} \frac{\pi^2}{30} g_*(T_{\gamma,m}) T_{N_1,m}^4$$

$$\rho_{N_2}(t_m) = m_{N_2} Y_{N_2}(z_m) \frac{2\pi^2}{45} g_*(T_{\gamma,m}) T_{\gamma,m}^3$$

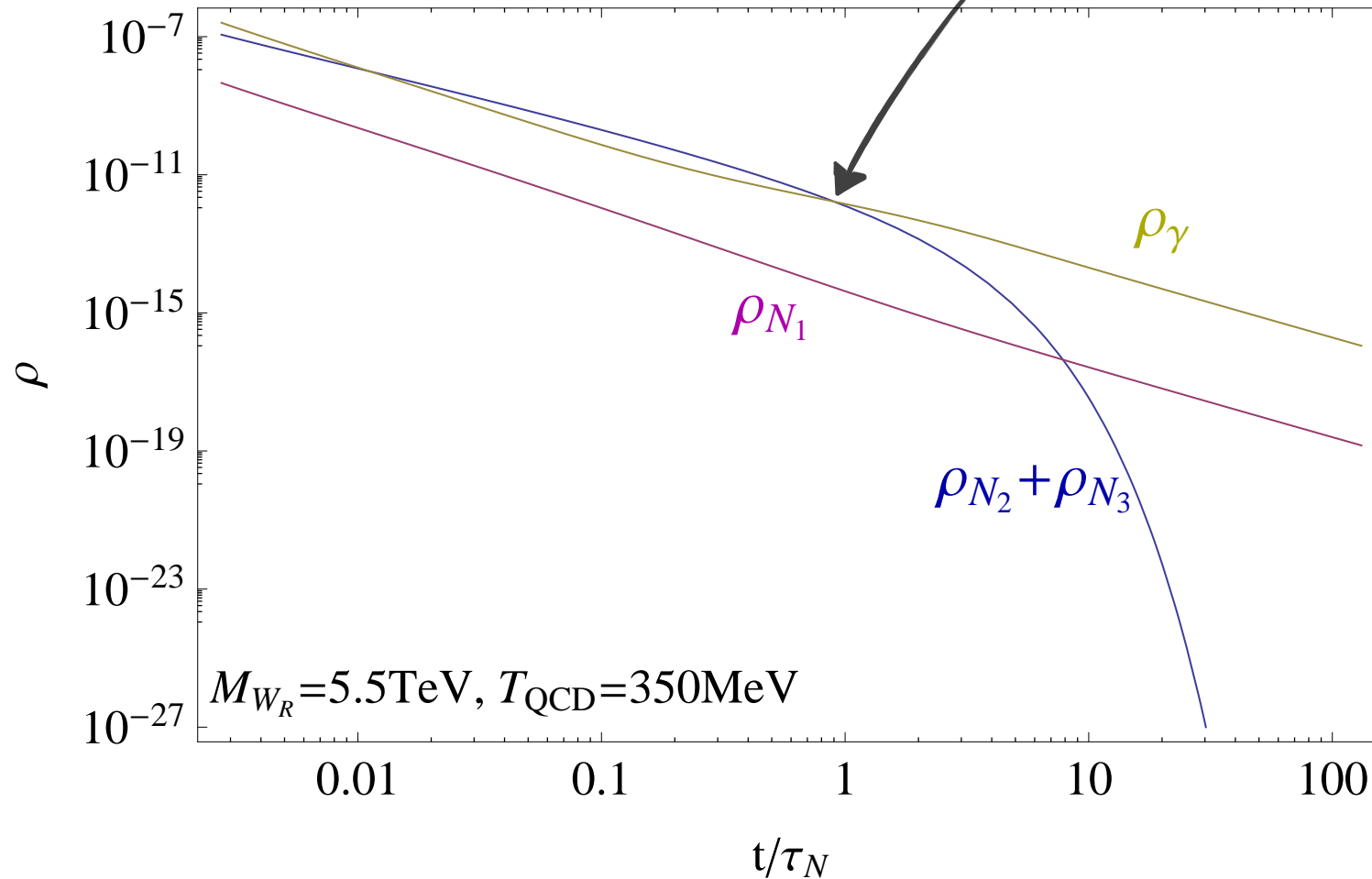
$$\rho_{N_3}(t_m) = m_{N_3} Y_{N_3}(z_m) \frac{2\pi^2}{45} g_*(T_{\gamma,m}) T_{\gamma,m}^3$$

# Dilution part II



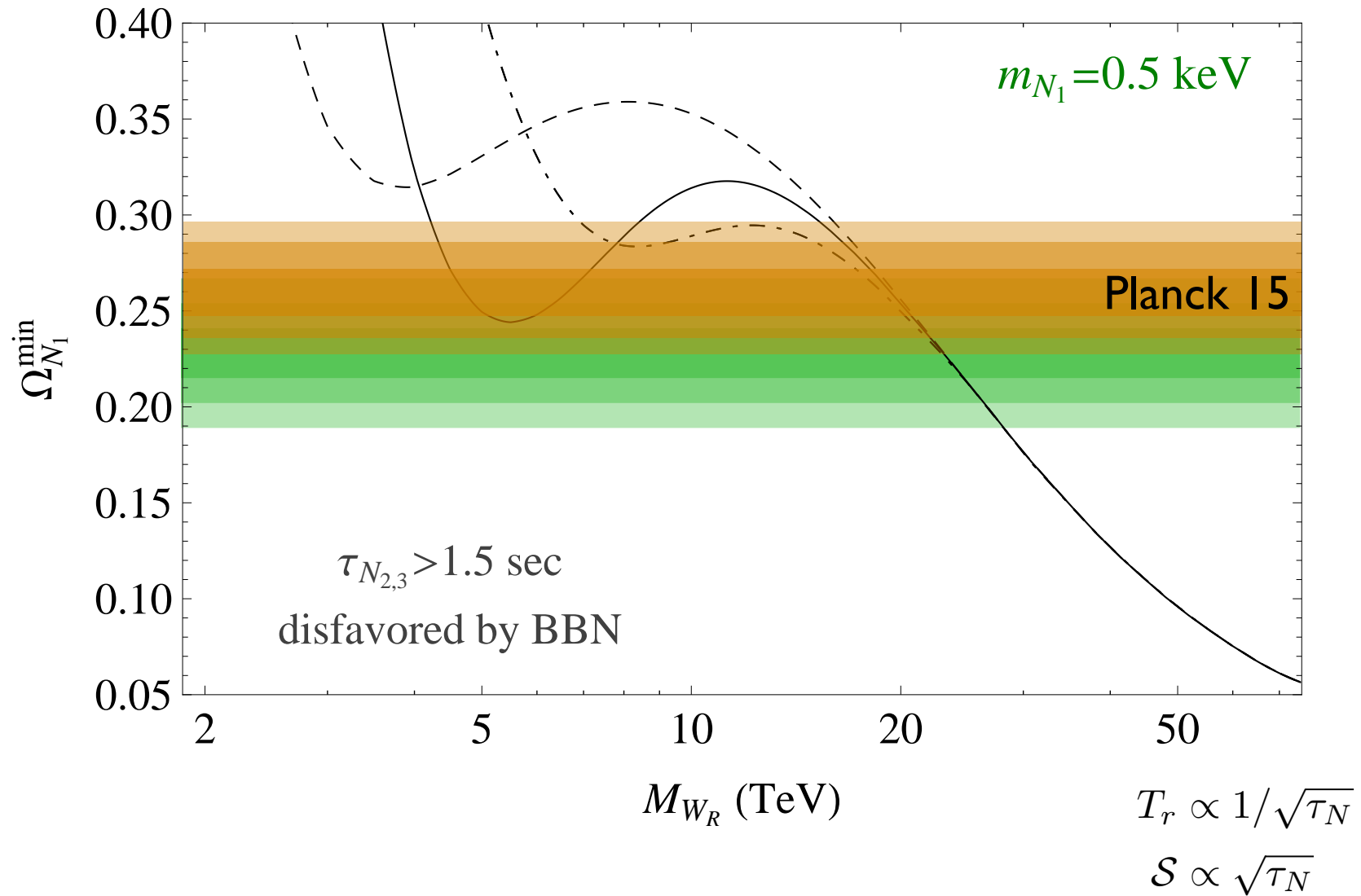
# Dilution part II

$$\rho_{N_2} + \rho_{N_3} = \rho_\gamma = \frac{\pi^2}{30} g_*(T_r) T_r^4 \quad T_r \gtrsim 0.7 \text{ MeV}$$



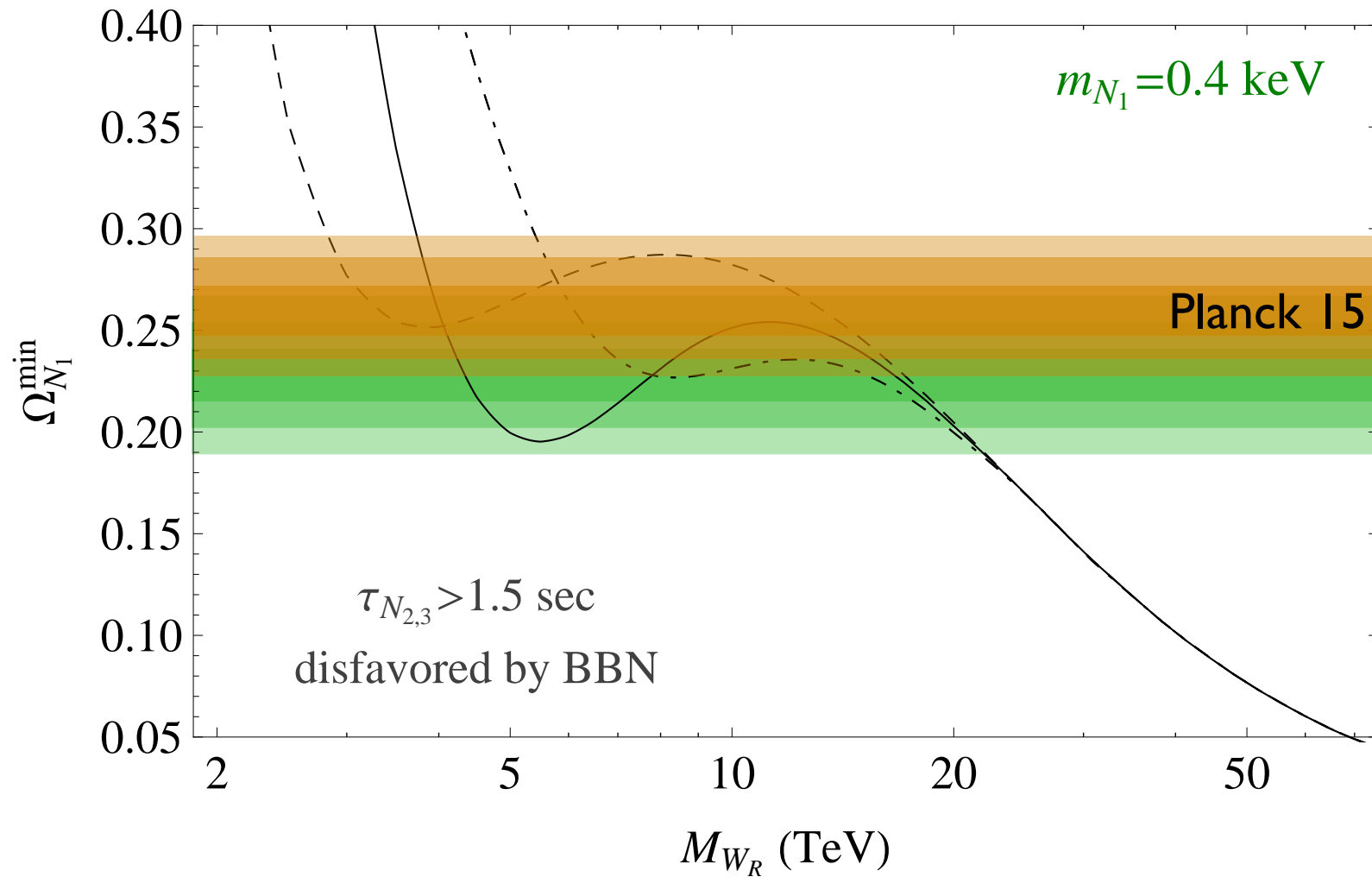
$$S = \frac{S(t_f)}{S(t_m)} = \frac{s(t_f)}{s(t_m)} \frac{V(t_f)}{V(t_m)} = \left( \frac{\rho_\gamma(t_f)}{\rho_\gamma(t_m)} \right)^{3/4} \left( \frac{\rho_{N_1}(t_f)}{\rho_{N_1}(t_m)} \right)^{3/4}$$

# Left-Right low scale DM windows





# Left-Right low scale DM windows



## Left-Right low scale DM - further constraints

— dSPH

$$m_{N_1} \lesssim (0.4 - 0.5) \text{ keV}$$

— diluted Ly- $\alpha$

$$m_{N_1} \lesssim \mathcal{O}(\text{keV})$$

— CMB &  $N_{eff}$  repopulation

$$N_3 \rightarrow \pi^+ e^- \rightarrow e^+ e^- \nu_\mu \bar{\nu}_\mu \nu_e$$

$$N_2 \rightarrow \pi^+ \mu^- \rightarrow e^+ e^- \nu_\mu \nu_\mu \bar{\nu}_\mu \nu_e \bar{\nu}_e$$

Fuller,  
Kishimoto,  
Kusenko '11

— X-ray observations

$$M_D, \xi_{LR} \text{ \& } v_L$$

Merle, Niro '13

—  $0\nu 2\beta$  subject to uncertainties  $M_{W_R} \gtrsim (5 - 7) \text{ TeV}$

MN, Senjanović, Zhang '12  
Lopez-Pavon, Huang '13

# Critical conclusion

pros

- RH neutrino keV DM candidate
- dilution possible in the minimal model
- window at the LHC

cons

- X-ray limits
- fairly low reheating vs. BBN and CMB
- light DM mass vs. Ly- $\alpha$

Thank you

# X-ray constraints

## Majorana vs. Dirac in Left-Right

MN, Senjanović, Tello '13

seesaw  $M_\nu = -M_D^T M_N^{-1} M_D + \frac{v_L}{v_R} M_N$

parity  $M_D = M_D^T$

$$M_N = V_R^T m_N V_R$$

$$M_D = i M_N \sqrt{M_N^{-1} M_\nu - \frac{v_L}{v_R}}$$

$$\mathcal{A}_{N_1 \rightarrow \nu \gamma} = \text{diagram}_1 + \text{diagram}_2$$

The diagram shows the amplitude  $\mathcal{A}_{N_1 \rightarrow \nu \gamma}$  as the sum of two terms. The first term is a tree-level diagram where an incoming  $N_1$  line splits into  $\nu_i$  and  $\nu_k$  lines, with a  $W$  boson loop connecting them, and a  $\gamma$  photon is emitted from the loop. The second term is a tree-level diagram where an incoming  $N_1$  line splits into  $\nu_j$  and  $\nu_i$  lines, with a  $\xi_{LR}$  loop connecting them, and a  $\gamma$  photon is emitted from the loop.

$$\Gamma_{N_1 \rightarrow \nu \gamma} = \Gamma_\mu \left( \frac{m_{N_1}}{m_\mu} \right)^5 \left( \frac{\alpha}{4\pi} \right)^3 \sum_\ell \left| \theta_{LR\ell N_1} f_D + \xi_{LR} \frac{m_\ell}{m_{N_1}} f_{\xi_{LR}} \right|^2$$

# $0\nu 2\beta$ constraints

$$m_{ee}^{N_3} = p^2 \left( \frac{M_W}{M_{W_R}} \right)^4 V_{Re3}^2 \frac{m_{N_3}}{p^2 + m_{N_3}^2}$$

MN, Nesti, Senjanović, Tello '11

including the mixing

$$m_{ee}^{\nu N} = \left( \xi_{LR} + \eta \frac{M_W^2}{M_{W_R}^2} \right) p (M_N^{-1} M_D)_{ee}$$

MN, Senjanović, Tello '12

see also

Lopez-Pavon, Huang '13

Barry, Rodejohann '13

NME uncertainties

axial coupling

Barea, Kotila, Iachello '12, '13, '15

